

# Good approximate QLDPC codes from spacetime Hamiltonians

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joint work with

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[arXiv:1811.00277](https://arxiv.org/abs/1811.00277)

QIP 2019

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# Why study error-correcting codes?

## Quantum fault tolerance

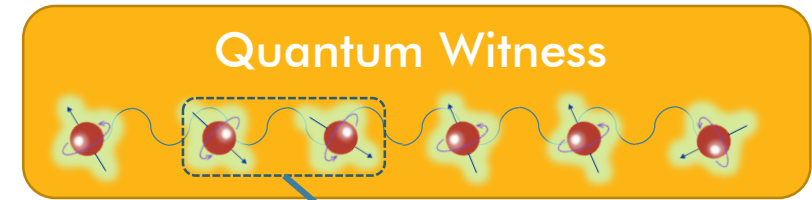
- [Gottesman<sup>09</sup>] QLDPC  $\Rightarrow$  fault tolerance quantum computation with constant overhead

## Quantum PCP conjecture

- Hardness of approximation in quantum setting
- Entanglement at room temperature

## Interesting local Hamiltonians

- with robust entanglement properties
- toric code, color codes, etc.



input  $x$

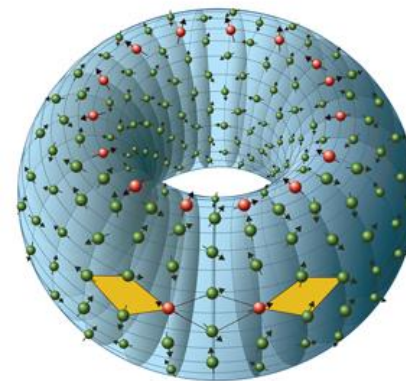
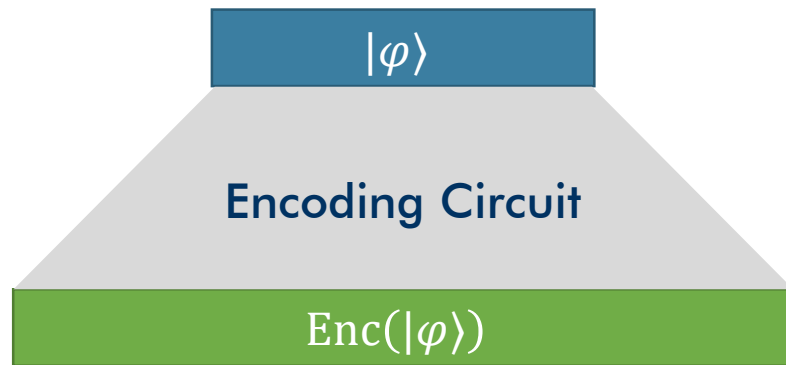


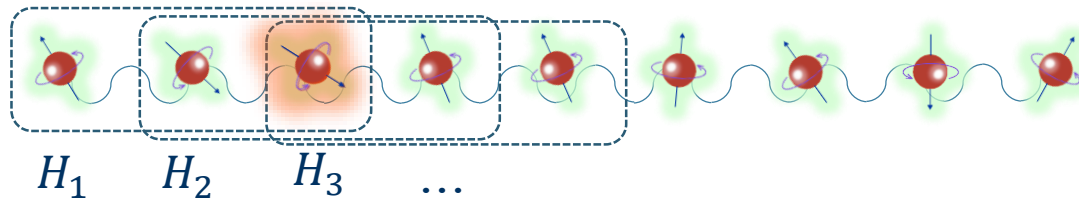
Image: Daniel Gottesman, APS

# What makes a code good?

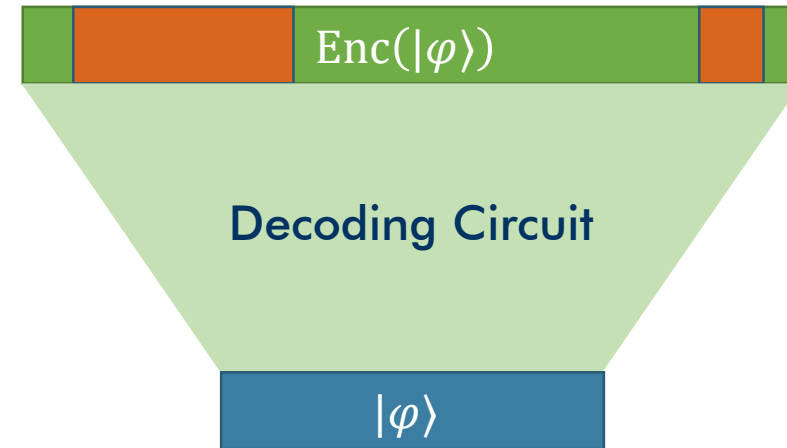
Rate



Stabilizer weight (locality)



Distance



$$\text{Rate: } \frac{k}{n} = \Omega(1)$$

$$\text{Distance: } d = \Omega(n)$$

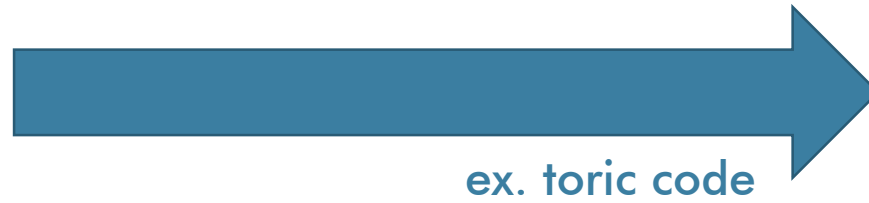
$$\text{Locality: } O(1)$$

We show that optimal rate, distance and locality parameters *are possible* (modulo polylog corrections)

if we go beyond stabilizer codes to

**non-commuting and approximate codes**

Quantum error  
correcting codes



ex. toric code

Local  
Hamiltonians

(with robust entanglement)



this talk

# Outline

- Coding theory definitions
- Uniformization via sorting circuits
- Spacetime Hamiltonians
- Spectral gap analysis

# What is a LDPC code?

Classically, a code  $\mathcal{C}$  is a dim  $k$  subspace of  $\mathbb{Z}_2^n$ .

A linear code can be defined by a matrix  $H \in \mathbb{Z}_2^{n \times (n-k)}$ .

$$\mathcal{C} = \{x \in \mathbb{Z}_2^n : Hx = 0\}$$

**L**ow  
**D**ensity  
**P**arity  
**C**heck

$$H = \begin{pmatrix} 1 & 1 & 0 \\ 0 & 1 & 1 \end{pmatrix} \quad \begin{pmatrix} 1 & 1 & 0 \\ 0 & 1 & 1 \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \\ x_3 \end{pmatrix} = 0$$

$$x_1 = x_2 = x_3 \quad \Rightarrow \quad \mathcal{C} = \{000, 111\}$$

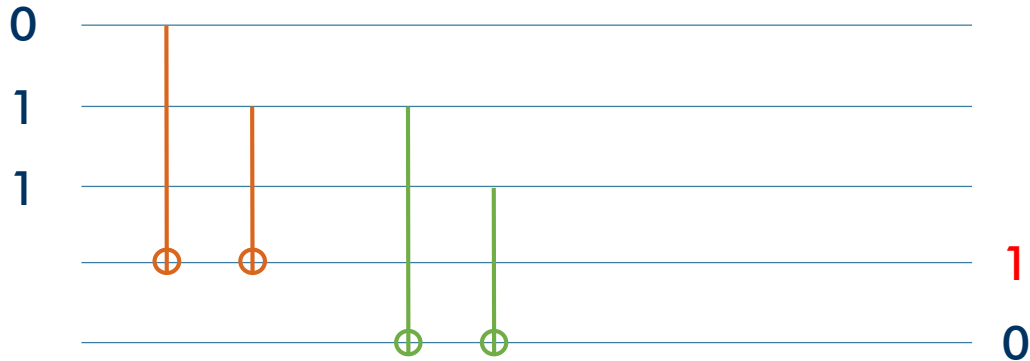
$H$  has  $c$ -locality if  $H$  is  $c$ -row sparse and  $c$ -column sparse.

# Benefits of an LDPC code

$$H = \begin{pmatrix} 1 & 1 & 0 \\ 0 & 1 & 1 \end{pmatrix}$$

Since the checks overlap, they can't be parallelized and must be done in series.

If the code is  $c$ -local, then the checks can be parallelized into  $c^3 + c$  depth circuit.



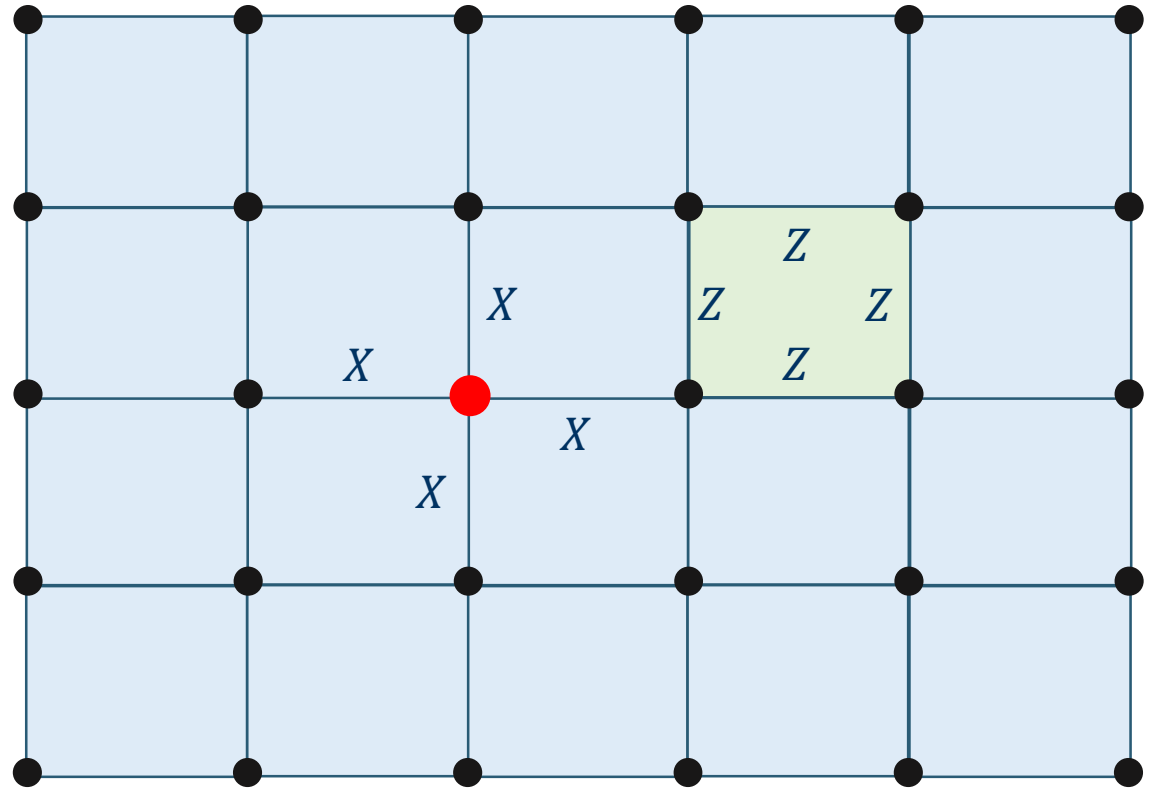
Proof: Each check shares bits with at most  $c^2$  other checks. By coloring argument, requires  $c^2 + 1$  rounds. Each round requires depth  $c$ .



# Quantum LDPC codes

For CSS codes  
(codes that handle  
 $X$  errors and  $Z$   
errors separately),  
definition is easy...

both parity check  
matrices  $H_X$  and  $H_Z$   
need to have low  
density.



ex. toric code

locality: 4

rate:  $2/n$

distance:  $O(\sqrt{n})$

Can we do better?

# Best known stabilizer codes

- [Tillich-Zemor<sup>13</sup>]
  - rate:  $\Omega(1)$
  - distance:  $O(\sqrt{n})$
  - locality:  $O(1)$
- [Freedman-Meyer-Luo<sup>02</sup>]
  - rate:  $\Omega(1/n)$
  - distance:  $O(\sqrt{n \log n})$
  - locality:  $O(1)$

To do better, we probably need to go past stabilizer codes!

# Going past stabilizer codes

Let  $H_1, H_2, \dots, H_m$  be a set of  $c$ -local projectors acting on  $n$  qubits. ← not necessarily commuting

Define the code-space  $\mathcal{C}$  as the mutual eigenspace:

$$\mathcal{C} = \{|\varphi\rangle \in (\mathbb{C}^2)^{\otimes n} \mid \langle \varphi | H_i | \varphi \rangle = 0 \ \forall H_i\}$$

$H = H_1 + \dots + H_m$  is  $c$ -QLDPC if additionally each qubit participates in at most  $c$  terms  $H_i$ .

# First attempt [N-Vazirani-Yuen<sup>18</sup>]

CSS codes exist with linear rate and distance, but lack locality.

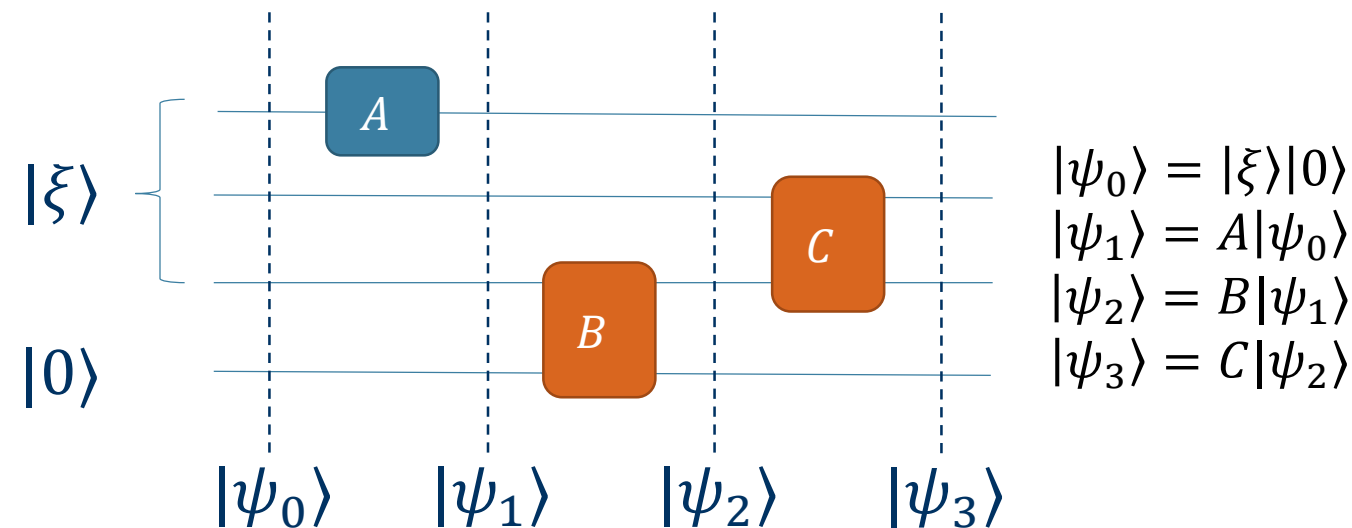


Create a Hamiltonian whose ground-space is almost exactly that of a CSS code but is locally checkable.

# First attempt [N-Vazirani-Yuen<sup>18</sup>]

Create a Hamiltonian whose ground-space is almost exactly that of a CSS code but is locally checkable.

Express a computation as the ground-state of a 5-local Hamiltonian (Feynman-Kitaev clock Hamiltonian) [Kitaev<sup>99</sup>]



Together,  $\{|\psi_t\rangle\}$  are a “proof” that the circuit was executed correctly. But,  $|\tilde{\Psi}\rangle = |\psi_0\rangle|\psi_1\rangle \dots |\psi_T\rangle$  is not locally-checkable.

Instead, the following “clock” state\* is:

$$|\Psi\rangle = \frac{1}{\sqrt{T+1}} \sum_{t=0}^T |t\rangle|\psi_t\rangle$$

\*Quantum analog of Cook<sup>71</sup>-Levin<sup>73</sup> Tableau.

# First attempt [N-Vazirani-Yuen<sup>18</sup>]

Create a Hamiltonian whose ground-space is almost exactly that of a CSS code but is locally checkable.

Let  $C = C_T C_{T-1} \dots C_1$  be a circuit with gates  $\{C_i\}$  and let  $|\psi_0\rangle = |\xi\rangle|0\rangle^{\otimes n-k}$  be an initial state for  $|\xi\rangle \in (\mathbb{C}^2)^{\otimes k}$ .

There is a local Hamiltonian with ground space of:

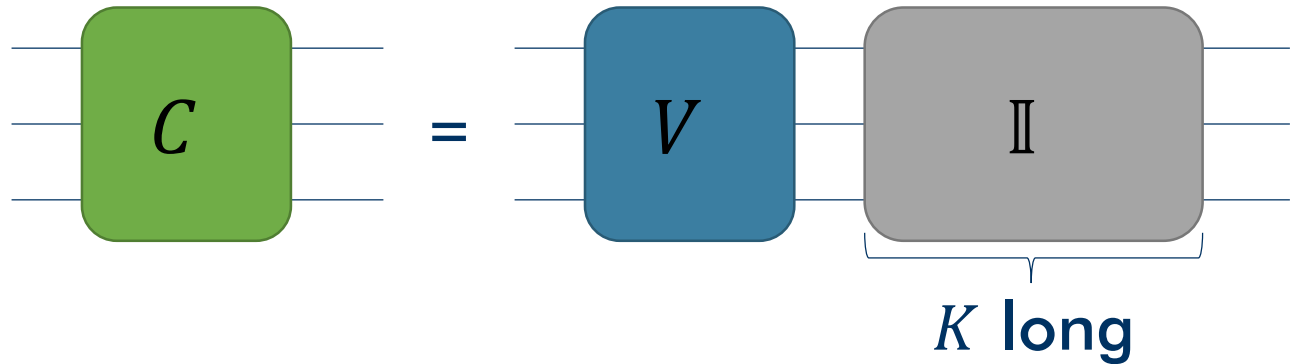
$$\mathcal{G} = \left\{ |\Psi_\xi\rangle = \frac{1}{\sqrt{T+1}} \sum_{t=0}^T |\text{unary}(t)\rangle \otimes |\psi_t\rangle : \begin{array}{l} |\psi_t\rangle = C_t |\psi_{t-1}\rangle, \\ |\psi_0\rangle = |\xi\rangle|0\rangle^{\otimes (n-k)} \end{array} \right\}.$$

# First attempt [N-Vazirani-Yuen<sup>18</sup>]

Create a Hamiltonian whose ground-space is almost exactly that of a CSS code but is locally checkable.

Let  $V$  be the encoding circuit for a good CSS code.

Choose  $K = O(T_V \delta^{-2})$ .



Construct the clock Hamiltonian for this “padded” circuit  $C$ .

# First attempt [N-Vazirani-Yuen<sup>18</sup>]

The groundspace of  $H$  is  $\approx$  the groundspace of a CSS code tensored with junk.

$$\mathcal{G}_C = \left\{ \frac{1}{\sqrt{T_C + 1}} \sum_{t=0}^T |t\rangle |\psi_t\rangle : \begin{array}{l} |\psi_t\rangle = C_t C_{t-1} \dots C_1 |\psi_0\rangle, \\ |\psi_0\rangle = |\xi\rangle |0\rangle^{\otimes (n-k)} \end{array} \right\}$$

But for  $t \geq T_V$ ,  $|\psi_t\rangle = V |\psi_0\rangle$ . Thus,  $1 - O(\delta^2)$  fraction of  $|\psi_t\rangle = V |\psi_0\rangle$ .

$$\mathcal{G}_C \approx \frac{1}{\sqrt{T_C + 1}} \sum_{t=0}^T |t\rangle \otimes \{V |\psi_0\rangle : |\psi_0\rangle = |\xi\rangle |0\rangle^{\otimes (n-k)}\}.$$

Plus,  $\mathcal{G}_C$  is the ground-space of a 5-local Hamiltonian!

However, some qubits participate in many terms  $H_t$ .



# First attempt [N-Vazirani-Yuen<sup>18</sup>]

However, some qubits participate in many terms  $H_t$ .

$H_t$  checks that the slice  $|t\rangle|\psi_t\rangle$  and the slice  $|t+1\rangle|\psi_{t+1}\rangle$  satisfy

$$|\psi_{t+1}\rangle = U_t|\psi_t\rangle$$

  $t^{\text{th}}$  gate of circuit

Locality of the code corresponds to the connectivity of the qubits in the circuit.

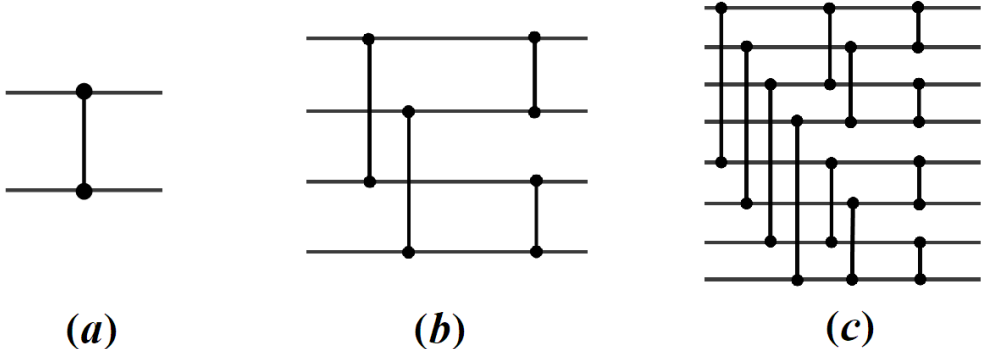


Minimize connectivity of the qubits in the circuit.

# Localizing the circuit via bitonic sorting circuits

Minimize connectivity of the qubits in the circuit.

Theorem [Batcher<sup>65</sup>]: There is a circuit of depth  $\log^2 n$  with  $\log n$  connectivity sorting  $n$  elements.

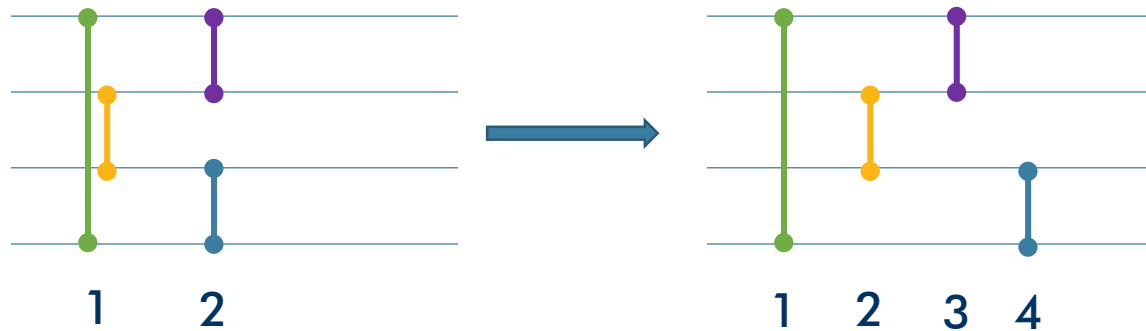


Can stretch circuit by  $\log^2 n$  mult. depth and reduce connectivity to  $n$ .



Can be used anywhere to simplify circuit connectivity in any situation.

# Long clocks and brittle Hamiltonians



For Feynman-Kitaev clock Hamiltonian each layer of the circuit needs exactly 1 gate.

This yields long clocks and brittle Hamiltonians.

**Brittle Hamiltonian:** Small spectral gap.

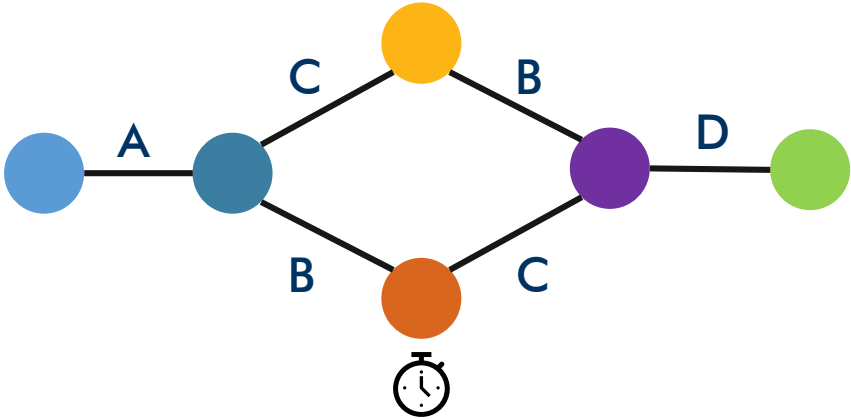
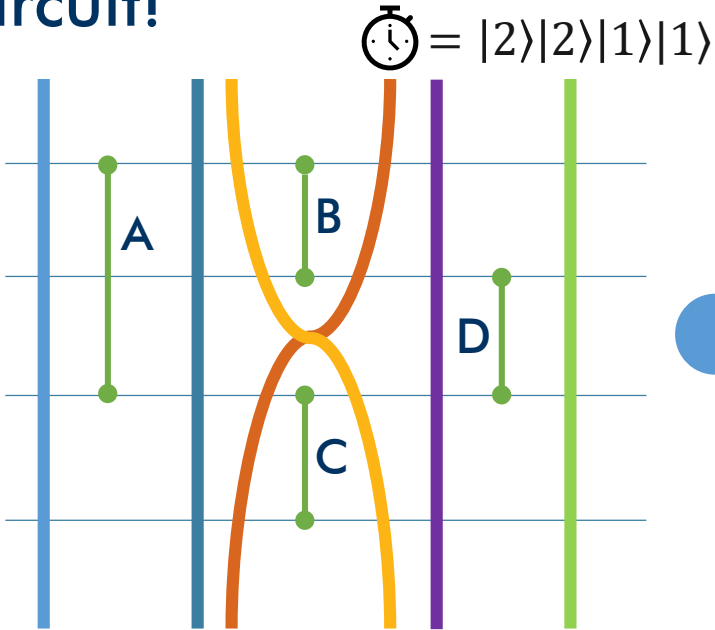
Not satisfying any 1 equation of  $|\psi_{t+1}\rangle = U_t|\psi_t\rangle$  has energy  $O(1/|C|)$  with  $|C| =$  the number of gates in  $C$ .

# Long clocks and brittle Hamiltonians

This yields long clocks and brittle Hamiltonians.



There are more than  $|C|$  partial computations of a circuit!



Build Hamiltonian with ground-state of uniform superposition over all partial computations  $\tau$ :

$$\sum |\tau\rangle|\psi_\tau\rangle$$

Space-time Hamiltonian  
[Breuckmann-Terhal<sup>14</sup>]

# Long clocks and brittle Hamiltonians



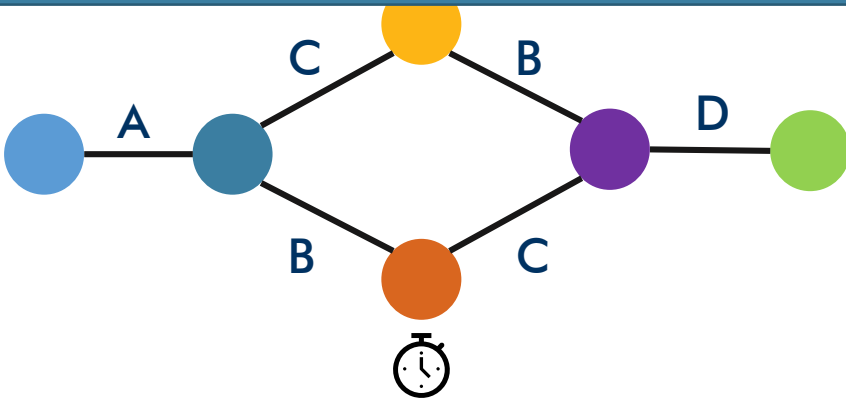
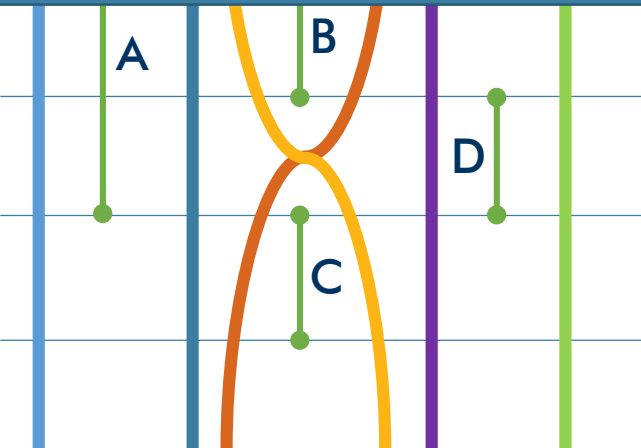
Theorem: This yields a Hamiltonian for whom the spectral gap scales

$$\tilde{O}\left(\frac{1}{n^{3.09} \text{depth}(C)^2}\right)$$

Instead of  $O\left(\frac{1}{|C|}\right)$  as in standard Feynman-Kitaev clock Hamiltonian

n with  
of  
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$$\sum |\tau\rangle |\psi_\tau\rangle$$



Space-time Hamiltonian [Breuckmann-Terhal<sup>14</sup>]

# Long clocks and brittle Hamiltonians



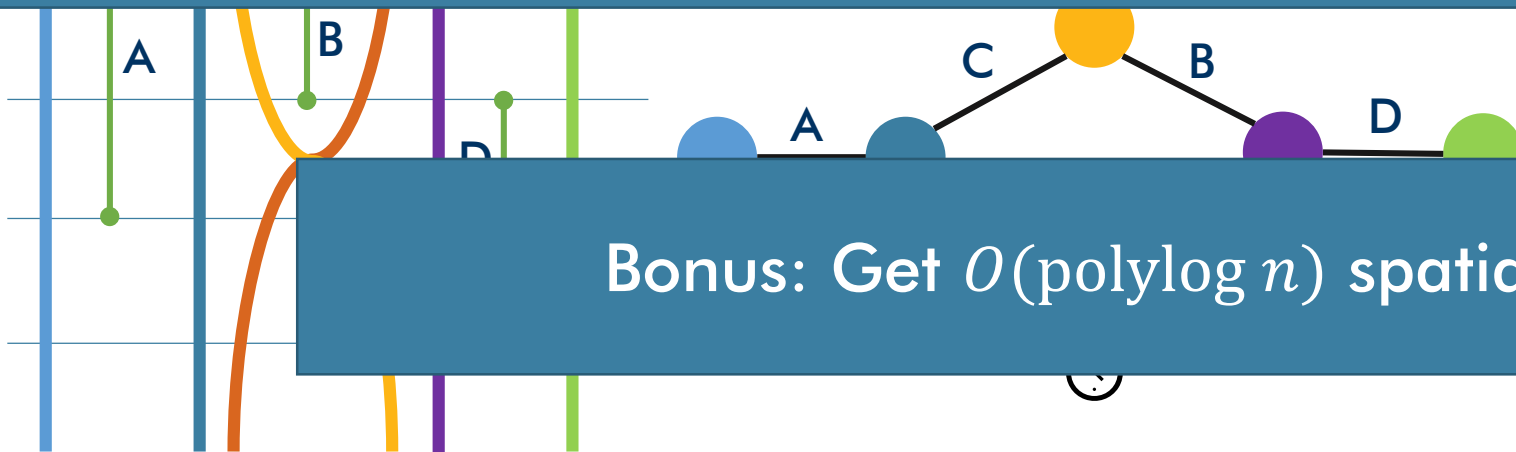
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Bonus: Get  $O(\text{polylog } n)$  spatial locality

[Breuckmann-Terhal<sup>14</sup>]

# Approximate decoding

$$|\Psi\rangle = \sum_{\tau} |\tau\rangle |\psi_t\rangle \approx \sum_{\tau} |\tau\rangle \otimes |\psi_{final}\rangle$$

↑  
over all valid partial  
computations

↑  
due to padding with  
identity gates

$$\mathcal{E}(|\Psi\rangle) \approx \mathcal{E}_1 \left( \sum_{\tau} |\tau\rangle \right) \otimes \mathcal{E}_2(|\psi_{final}\rangle)$$

Approximate decoding:

1. Trace out clock registers
2. Apply underlying code decoding procedure

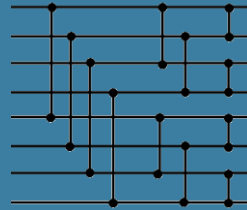
# Construction recap

A code with linear rate and distance and  $O(\log^3 n)$  depth encoding circuit [Brown-Fawzi<sup>13</sup>]

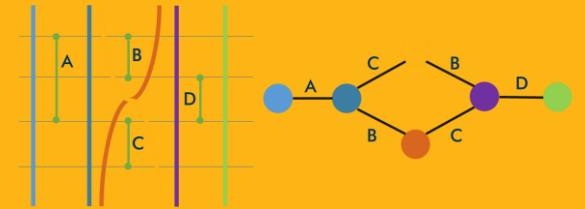
Pad the circuit with identity gates



Uniformize the connectivity of the circuit using bitonic sorting circuits



Build spacetime Hamiltonian of resulting code [Breuckmann-Terhal<sup>14</sup>]

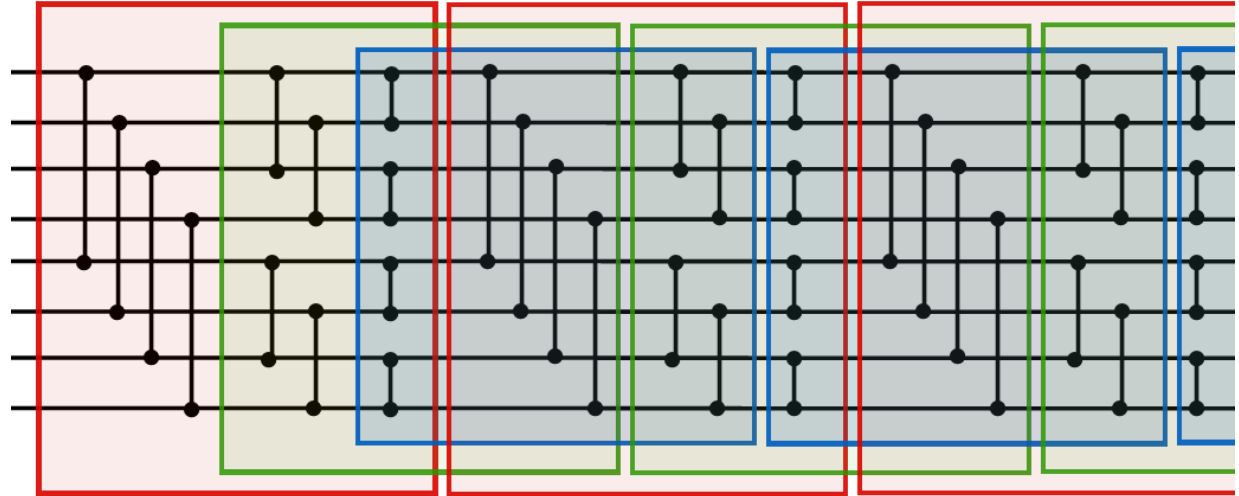
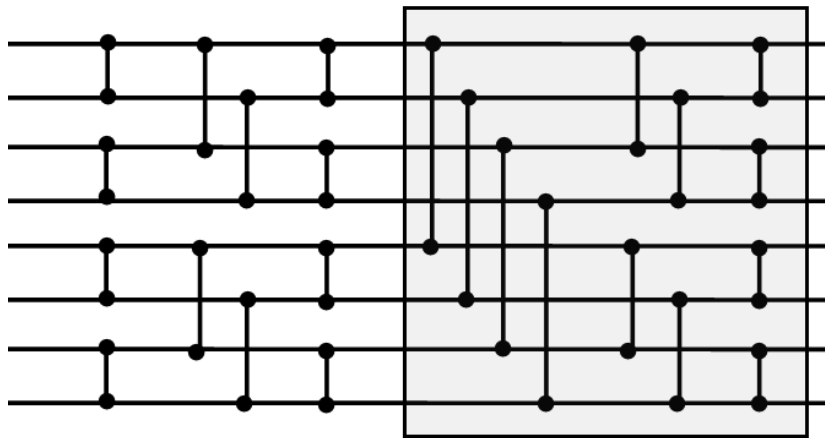




# Spectral gap analysis

Def: minimum non-zero eigenvalue of Hamiltonian  $H$

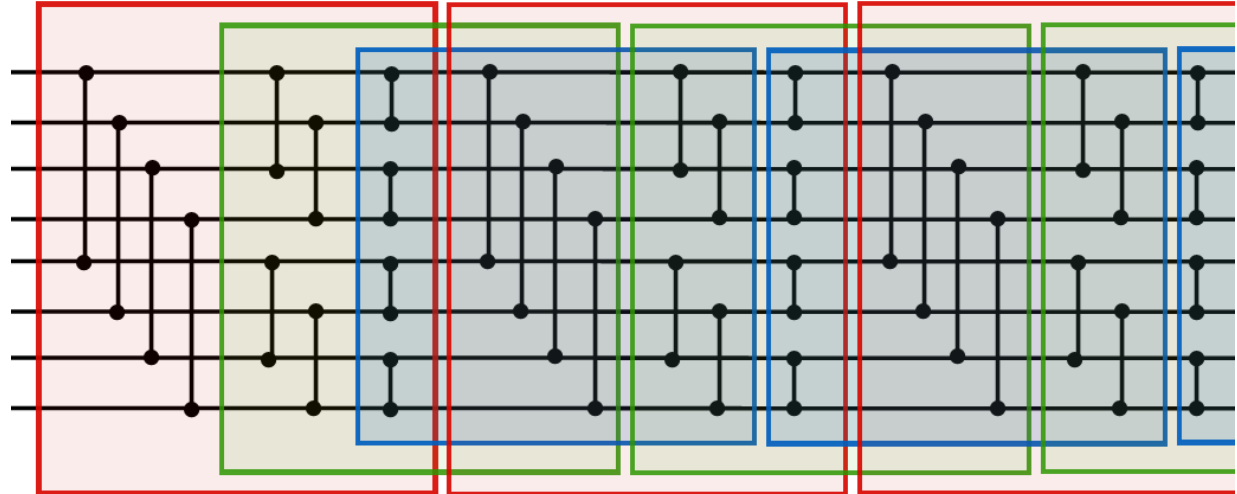
Map the Hamiltonian to a Markov chain over the space of valid partial computations



# Spectral gap analysis

Spectral gap of the code is based on the mixing time of valid configurations of a bitonic block

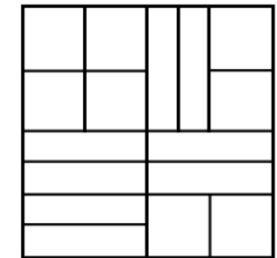
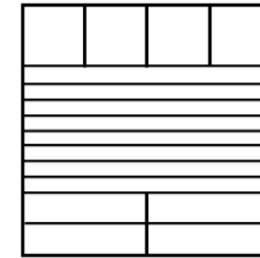
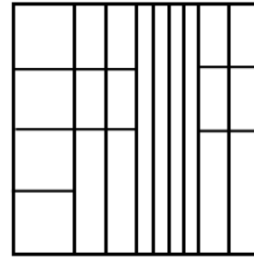
True of all constructions built from bitonic sorting circuits



# Spectral gap analysis

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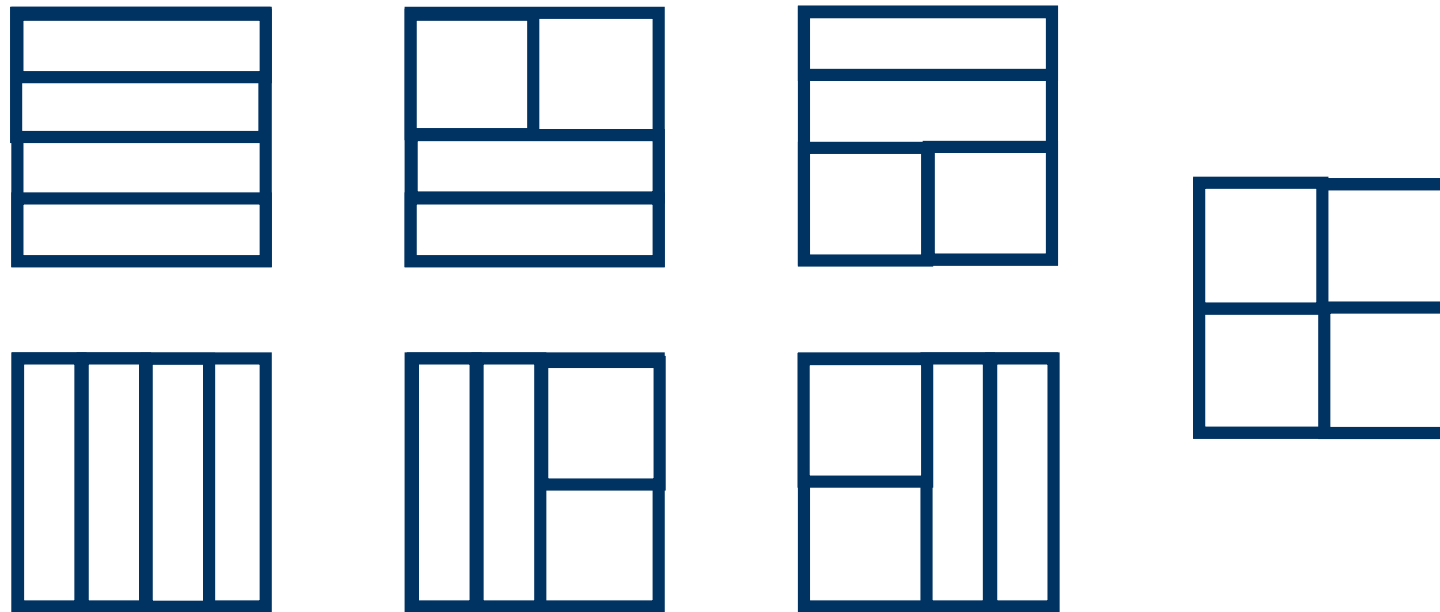
We noticed that bitonic blocks look similar to a structure called dyadic tilings studied in [Cannon-Levin-Stauffer<sup>17</sup>]



Dyadic tilings are ways of covering the unit square by  $2^d$  rectangles with corner coordinates at multiples of  $2^{-d}$

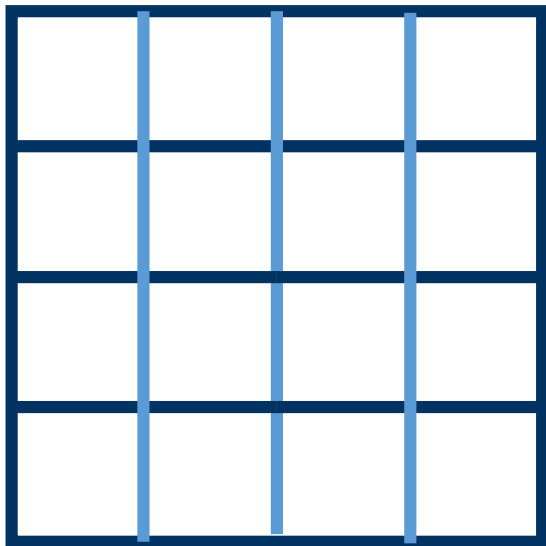
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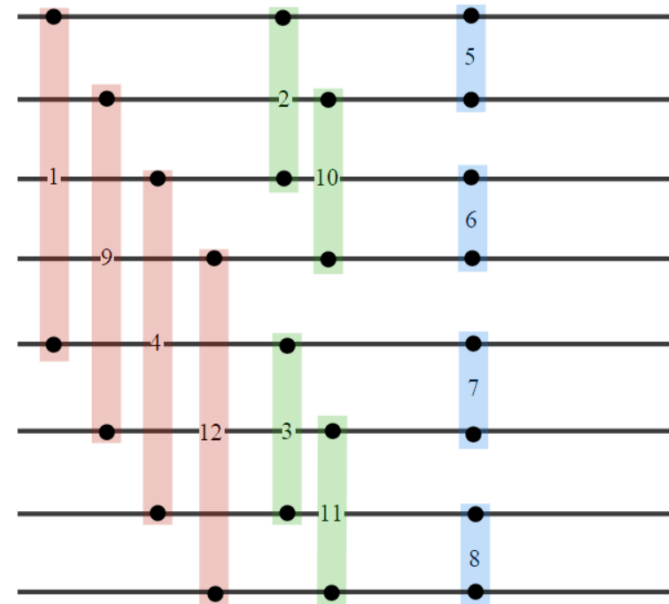
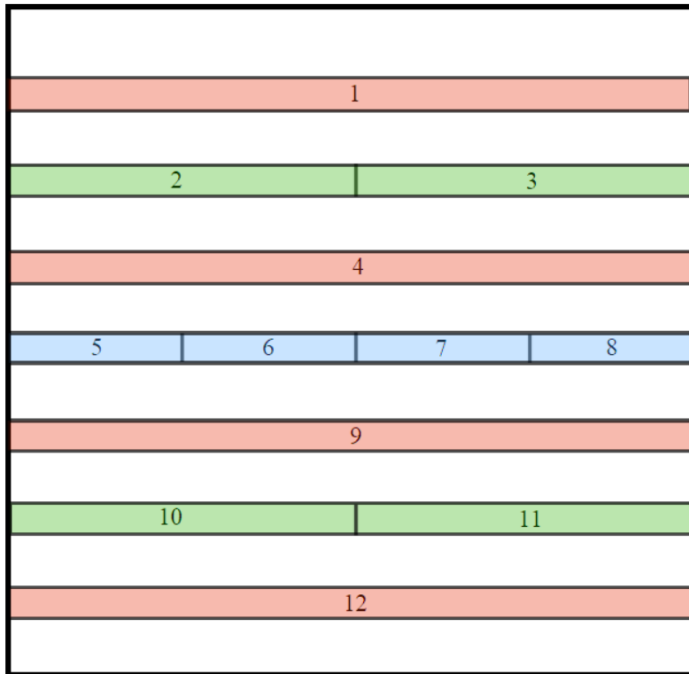
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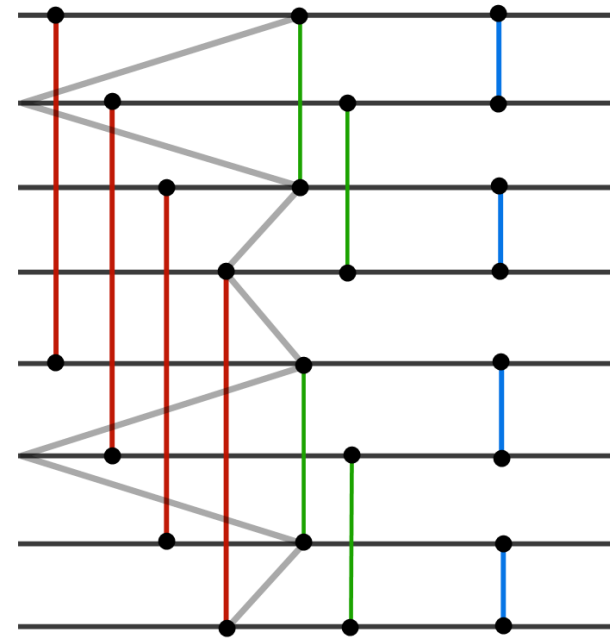
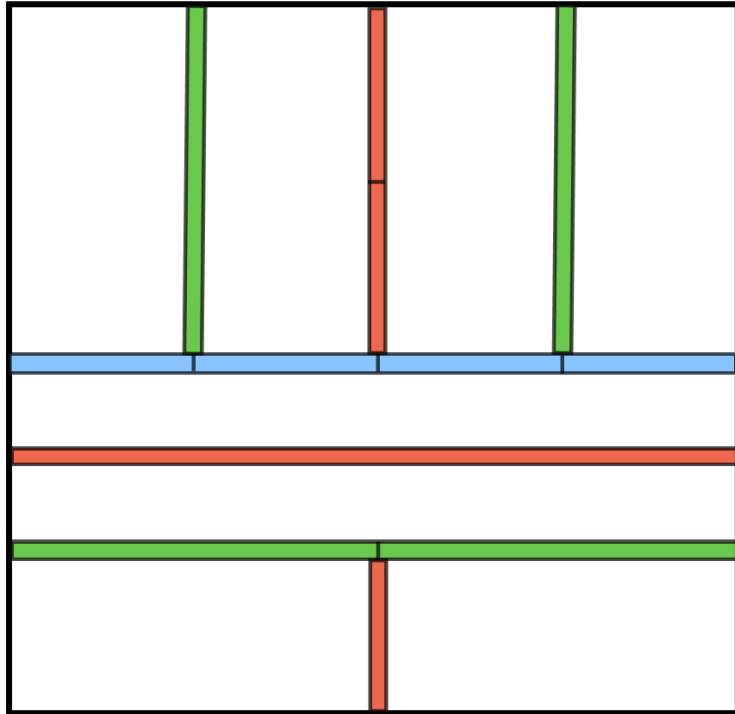
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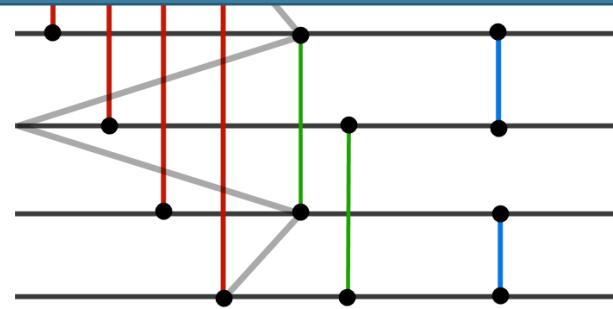
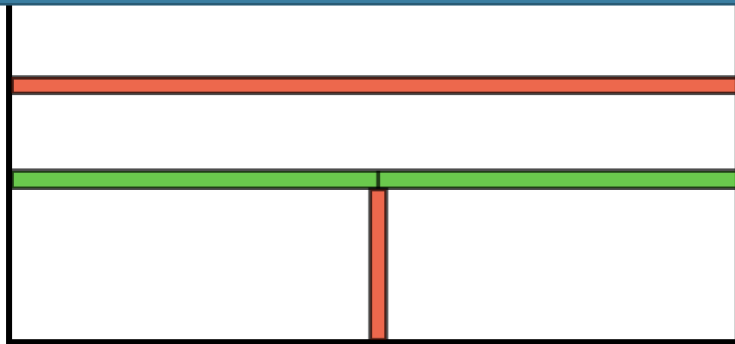


# Spectral gap analysis

Spectral gap of the code is based on the mixing time of valid configurations of a bitonic block

Theorem: The spectral gap of this Hamiltonian is

$$\tilde{\Omega} (n^{-3.09}).$$





# Summary of results

We constructed a new type of code based on spacetime Hamiltonians.

It has the following properties:

- rate:  $\Omega\left(\frac{1}{\text{polylog } n}\right)$
- distance:  $\Omega\left(\frac{n}{\text{polylog } n}\right)$
- spatial-locality:  $\Omega(\text{polylog } n)$
- spectral-gap:  $\Omega(n^{-3.09})$

Along the way, we also learned about

- localizing large stabilizers using circuit-to-Hamiltonian constructions
- uniformizing circuits with bitonic sorting networks
- analysis of uniform circuits via Markov chain techniques

# What does this teach us?

First, this isn't the "perfect" error-correcting code or is realistic

Relaxing the requirements of stabilizer codes is helpful

- Code-space as the ground-space of a sum of non-commuting projectors
- Approximate error-correction

There are connections between computation and error-correction that we don't fully understand!

# Thanks!