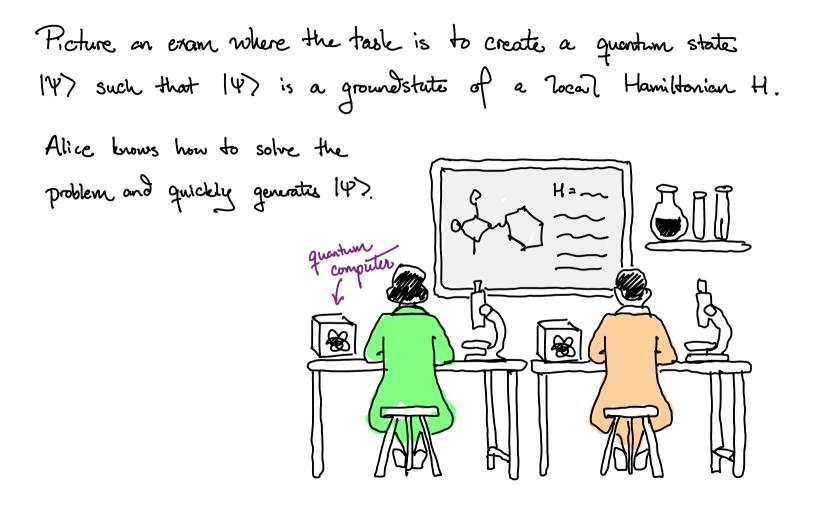
Based on a joint nork wich Vojtěch Harlíček

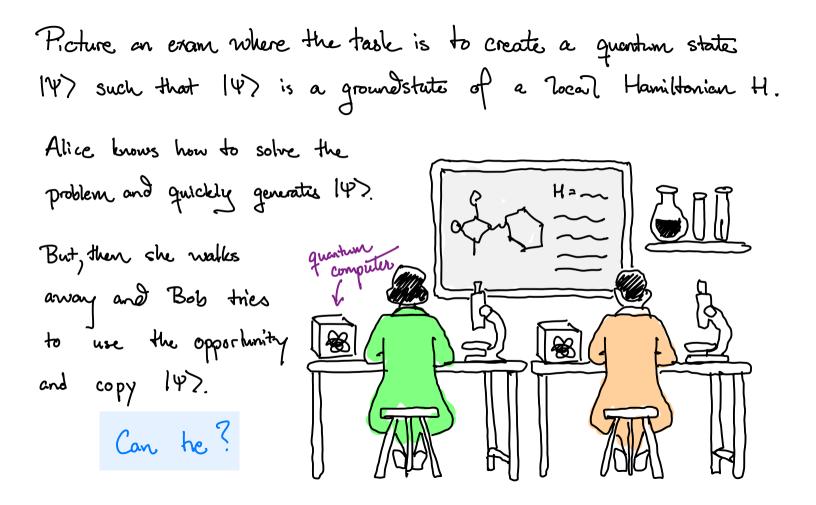
.

Picture on exam where the task is to create a quantum state 14> such that 14> is a groundstate of a local Hamiltonian H. Picture on exam where the task is to create a quantum states 14> such that 14> is a groundstate of a local Hamiltonian H.

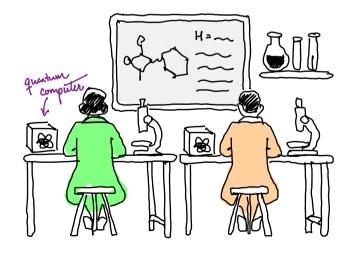




Picture on exam where the task is to create a quantum state 147 such that 147 is a groundstate of a rocar Hamiltonian H. Alice knows how to solve the problem and quickly generates 14>. But, then she walks away and Bob tries to use the opportunity and copy 142.







Carreats :

1) Bob's cheating must be efficient. He can't spend too long using 14>Alive. He can use the q. competters.



Carrents :

5

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## A first answer night trivially be <u>NO</u>.

There is no quantum transformation  
mapping 
$$|\Psi\rangle|0,...,0\rangle \mapsto |\Psi\rangle|\Psi\rangle$$

for all 147.

there is uncertainty about what (4) is.

If Bob knew nothing about the exam, he wouldn't be very good at cheating.

Either Alice or the Examiner will detect the malfiescance.

## But Bob has more information. In particular, he knows the exam question H.

<u>A notion of efficient copying</u>

A notion of efficient copying (time involvent)  
The Hamiltonion 
$$H = \sum h_i$$
 consists of terms acting on  
k out of n qubits. (Think k=2 or 3).  
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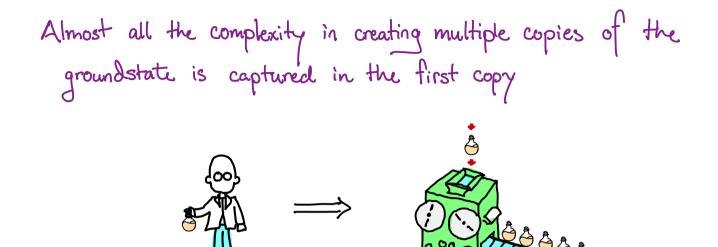
- solutions to q. problems can easily be publically disseminated - quantum copy-protection schemes can be broken

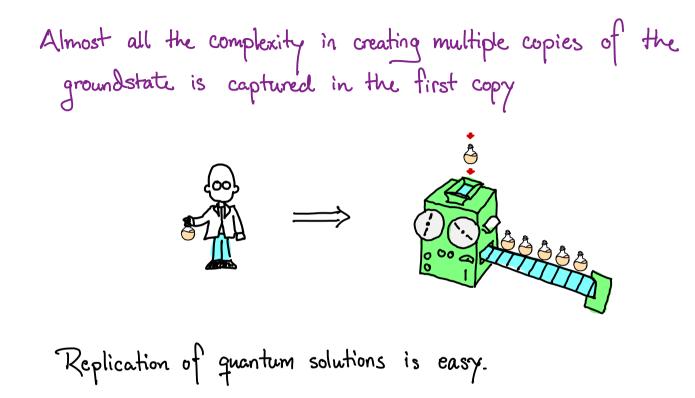
- solutions to q. problems can easily be publically disseminated
- quantum copy-protection schemes can be booleen
- quantum money can easily be counterfeited.

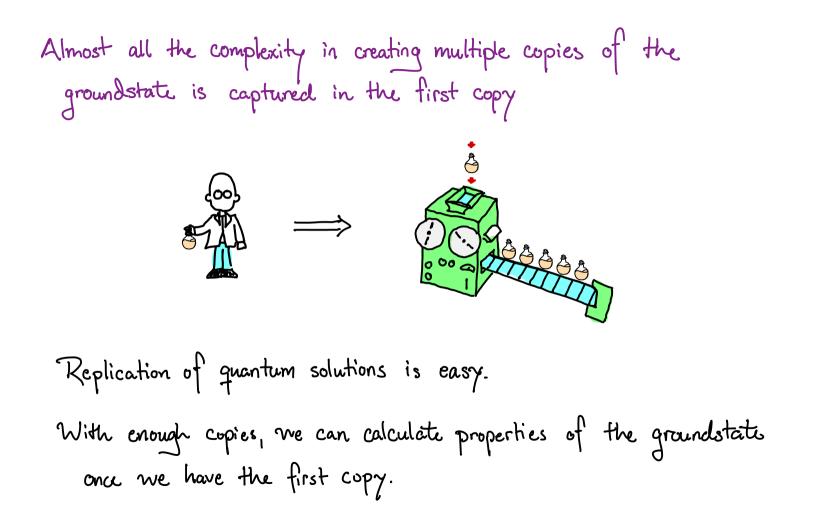
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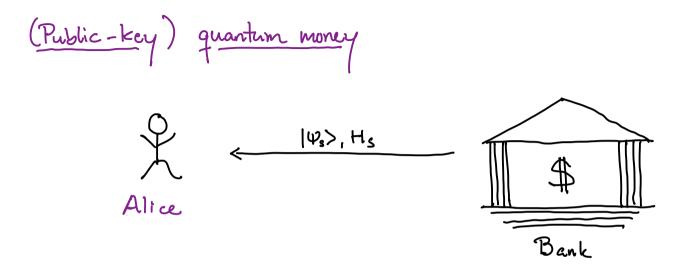


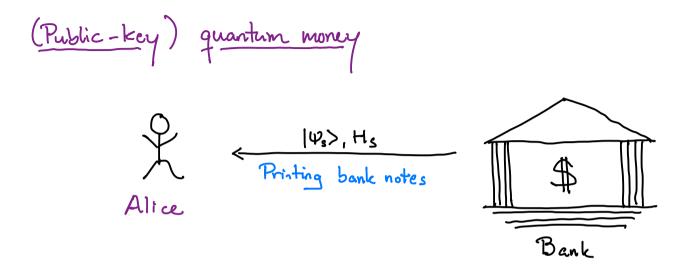
(Public-key) quarter money

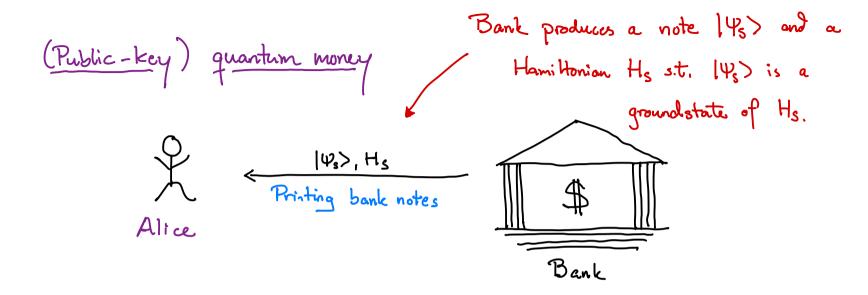
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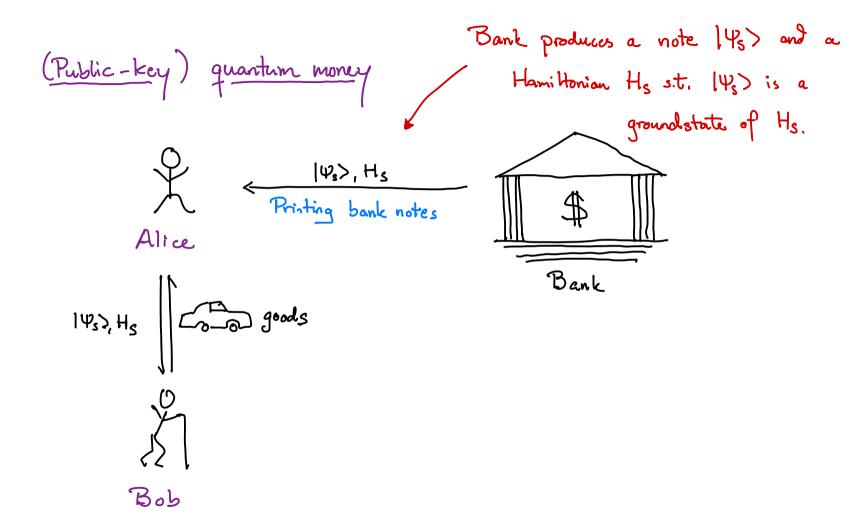
Alice

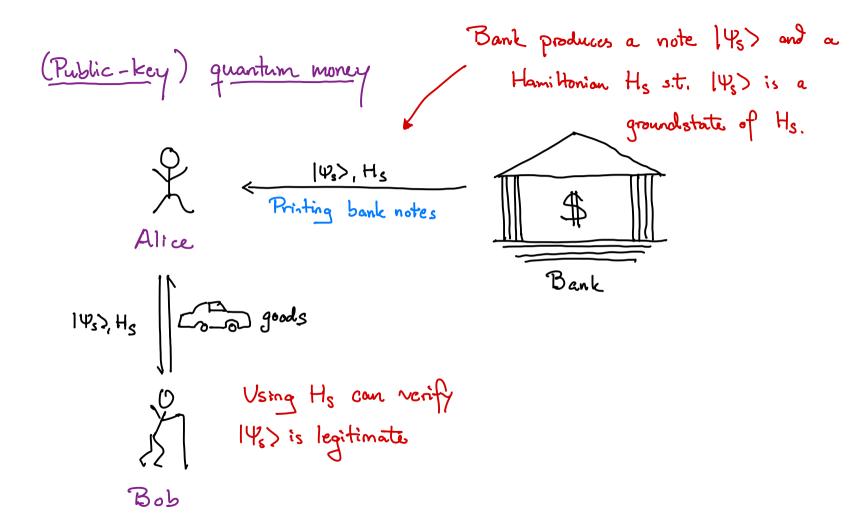














From a vague goal to an achievable mathematical theorem  
Proving outright a statement like "\_\_\_\_\_\_ transformation takes exponential time"  
is super-challenging. Akin to proving 
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The maximally entangled state for a subspace.

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The maximally entangled state for a subspace.

Let 
$$TT \subseteq (\overset{d_2}{\square})^{d_2}$$
 be a subspace defined by basis  $|b_1\rangle, ..., |b_d_1\rangle$ .  
The  $TT$ -maximally entangled state  $\in (\overset{d_2}{\square} \otimes (\overset{\alpha_2}{\square})^{d_2})^{d_2}$  is the state  $|\overline{\Phi}_T\rangle := \frac{1}{\sqrt{d_1}} \sum_{i=1}^{d_1} |b_i\rangle \otimes |b_i^*\rangle$ .

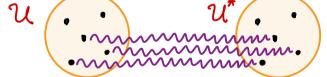
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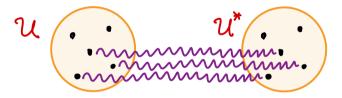
U

The TT-maximally entangled state  $\in ( \overset{d_2}{\circ} \otimes ( \overset{d_2}{\circ} ) )$  is the state  $| \overline{\Phi}_{TT} \rangle := \frac{1}{\sqrt{d_1}} \sum_{i=1}^{d_1} |b_i\rangle \otimes |b_i^*\rangle$ 



- generalizes EPR states to states over hidden subspaces.

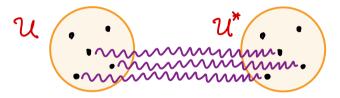
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which requires knowing how to map  $U: \mathbb{C}^{a_1} \hookrightarrow \mathbb{T}$ .

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   traition only!

Making intuitions formal

Conjecture:

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<u>Conjecture</u>:

# () Let V be an efficient quantum algorithm that accepts states in TT.

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(2) If I an algorithm W  

$$|\langle v \rangle \rangle |\overline{\Psi}_{\pi} \rangle |0\rangle \mapsto |\langle v \rangle \rangle |\overline{\Psi}_{\pi} \rangle |\overline{\Psi}_{\pi} \rangle$$

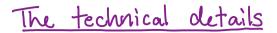
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Are states like 
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Are states like 
$$(\underline{\Psi}_{\mathrm{TT}})$$
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recall we wanted to consider physically relevant states.

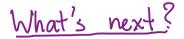


### The technical details

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which can be encoded in TT, H.

What's next?



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# What's next?

- Elephant in the noom: prove the conjecture previously stated
- -Discover other states which are hard to clove and find corresponding Hamiltonians - Prove hardness of cloning from state of the art assumptions like the learning with errors problem This work is definitively unfinished, but that's what makes it so tantalizing and exciting!