## The parametrized complexity of quantum verification <br> Srinivasan Arunachalam (IBM) <br> Sergey Bravyi (IBM) <br> Chinmay Nirkhe (UC Berkeley $\rightarrow$ IBM) <br> Bryan O'Gorman (IBM)

## The toggle between P and BQP



Counting the number of non-Clifford gates in a circuit is a measure of how "non-classical" a given circuit is

Gottesman \& Knill showed that there is a $P$ algorithm for deciding a quantum circuit decision problem if the circuit only has Clifford gates

A series of works has extended this famous theorem to Clifford gates of low non-Clifford gate count (most often counting the number of T gates) in both the decision and sampling regime

This gives parametrized (in \# of T gates $=t$ ) algorithms for quantum circuit problems.

## Is there a toggle between P and QMA?



Canonical QMA: Does there exist a $|\psi\rangle$ such that the circuit accepts with probability $>2 / 3$ or for all $|\psi\rangle$, is the acceptance probability bounded by $<1 / 3$ ?

What is the complexity of the parametrized quantum circuit satisfiability problem when the circuit on

- $n$ qubits,
- mancilla,
- $s$ gates,
- $t$ non-Clifford gates (T gates)?


## Is there a toggle between P and QMA?



## Non-determinism and quantum don't clash

- Yoganathan, Jozsa, and Strelchuk 2019 construct a reduction that reduces the computation (after classical processing) to a new $t$ T-gate computation on $n$ qubits witness but with no ancilla.
- In our result we maintain the ancilla but drastically reduces the witness to $t$ size.
- Furthermore, we give an $2^{\max (2+\alpha), \omega) \cdot t} \cdot \operatorname{poly}(s) \approx 5.3^{t}$ runtime algorithm for solving parametrized QCSAT
- $\alpha$ is stabilizer rank of magic states
- $\omega$ is matrix multiplication constant


## The Clifford perspective

- Clifford group $=\operatorname{span}(H, C N O T, S)$
- Classically simulable because $C P C^{\dagger}=P^{\prime}$ for any Pauli $P$

Warmup: Clifford QCSAT is in $P$
The measurement at the end $|0\rangle\langle 0|=\frac{I+Z}{2}$ so we are trying to optimize $\left.\frac{1}{2}+\frac{1}{2}\left|\langle\psi| \otimes\left\langle 0^{m}\right| C^{\dagger} Z C\right| \psi\right\rangle\left.\otimes\left|0^{m}\right\rangle\right|^{2}=\frac{1}{2}+\frac{\left.1\langle\psi| P|\psi\rangle \otimes\left\langle 0^{m}\right| Q\left|0^{m}\right\rangle\right|^{2}}{2}$ for Paulis $P, Q$.


Since we are trying to maximize $|\psi\rangle$ then $\langle\psi| P|\psi\rangle=1$ in best case.
3 cases: $\left.\frac{1}{2}+\frac{1}{2}\left|\left\langle 0^{m}\right| Q\right| 0^{m}\right\rangle\left.\right|^{2} \in\{0,1,1 / 2\}$. Can easily calculate given $Q$ and $Q$ is calculable using standard Clifford calculus.

## What happens when there are T gates

While Clifford conjugation maintain Paulis ...

$$
\begin{gathered}
T^{\dagger} I T=I \\
T^{\dagger} X T=(X+Y) / \sqrt{2} \\
T^{\dagger} Y T=(X-Y) / \sqrt{2} \\
T^{\dagger} Z T=Z
\end{gathered}
$$

So, by induction on the gates of a circuit $C$, we see that $C^{\dagger} Z C$ can be expressed as the lin. combination of $\leq 2^{t}$ terms

## Many terms but few linearly independent ones

Claim: There exists a basis of $t+1$ Pauli terms so that all $2^{t}$ terms can be expressed as products of basis terms.

Proof by induction: Base case: $I^{\dagger} Z I=Z=\mathrm{b}_{1}$.
Let $C=g_{s} g_{s-1} \ldots g_{1}$. At step $i$, let basis be $b_{1}, \ldots, b_{j}$.

- If $g_{i}$ is Clifford then, new basis of $g_{i}^{\dagger} b_{1} g_{i}, \ldots, g_{i}^{\dagger} b_{j} g_{i}$.
- If $g_{i}=T$ acting on qubit $q$,
- then first rewrite basis so that only $b_{1}, b_{2}$ act non-trivially on qubit $q$
- at most one of $b_{1}(q)$ and $b_{2}(q)$ is $\in\{X, Y\}$ and the other is $\{I, Y\}$
- wlog assume $b_{1}(q)=X$. Then add $b_{j+1}=b_{1} \cdot(X Y)_{q}$ to the basis.


## Many terms but few linearly independent ones

Claim: Since the $\leq 2^{t}$ Paulis have a linearly independent basis of $t+1$ Pauli terms, then there exists a Clifford unitary $W$ mapping these Pauli to a space of at most $t+1$ qubits.

Proof sketch: Each linearly independent Pauli defines a "qubit" and so $W$ can be constructed by a sequence of Clifford SWAP gate-like gadgets.

A more sophisticated analysis produces $W$ exactly with only poly ( $s$ ) pre-processing (not included in this talk).

## Interesting lower bound from upper bounds

- A reduction to a witness of length $t$
- A $5.3^{t} \cdot \operatorname{poly}(s)$ algorithm for solving parametrized QCSAT

Classical Exponential Time Hypothesis: SAT formulas on $n$ variables cannot be solved in time $2^{o(n)}$.

Corollary: There does not exist a generic reduction from SAT formulas on $n$ variables to SAT formulas on $o(n)$ variables.

## Interesting lower bound from upper bounds

- A reduction to a witness of length $t$
- A $5.3^{t} \cdot \operatorname{poly}(s)$ algorithm for solving parametrized QCSAT

ETH $\Rightarrow$ Quantum proof length optimality Conjecture: There does not exist a generic reduction from QCSAT formulas with witness length $n$ to QCSAT formulas with witness length $o(n)$.

Corollary: Assuming conj., in the worst case for QMA-hard problems, $t=\Omega(n)$.

## Lower bound for the complexity of $|W\rangle$

Any local Hamiltonian $H$ with $m$ terms can be expressed as the sum of $O(m)$ local Pauli terms.

Then there exists a Clifford operator $C$ s.t. $H=\left\langle W_{O(m)}\right| C\left|W_{O(m)}\right\rangle$
So, the local Hamiltonian problem can be expressed as

$$
\max _{\psi}\langle\psi, W| C|\psi, W\rangle
$$

Assume $|W\rangle=V\left|0^{k}\right\rangle$ for $V$ a circuit consisting of $t$ T-gates.

## Lower bound for the complexity of $|W\rangle$

Assume $\left|W_{O(m)}\right\rangle=V\left|0^{k}\right\rangle$ for $V$ a circuit consisting of $t$ T-gates.
Then, the problem can be rewritten as

$$
\max _{\psi}\langle\psi, W| C|\psi, W\rangle=\max _{\psi}\left\langle\psi, 0^{k}\right| V^{\dagger} C V\left|\psi, 0^{k}\right\rangle
$$

which can be reduced to (by main result) to witness length $t$.

Assuming optimal proof length conjecture, $t=\Omega(m)$ proving a linear lower bound on T-gate complexity of $|W\rangle$ state. Proof is robust to $1 / \operatorname{poly}(m)$ noise or $O(1)$ noise assuming QPCP.

## What's next

- A computational method for "testing" avg-case QMA vs QCMA
- Many other QMA-complete problems are built from q. circuits
- How many of them also have parametrized complexity solutions
- Ex. Non-identity check problem is in P for Clifford unitaries
- Is there a parameter like non-Clifford gate count that parametrizes the complexity of the local Hamiltonian problem?
- Is there a parameter that scales the problem between NP and QMA?
- What about between QCMA and QMA?

