

Quantum search-to-decision  
and the state synthesis problem

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joint work with

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How does one describe a quantum state?

How does one use a description of a quantum state?

Do quantum problems of classical description length  $l$  have classical solutions of length  $\text{poly}(l)$ ? (BCMA vs QMA)

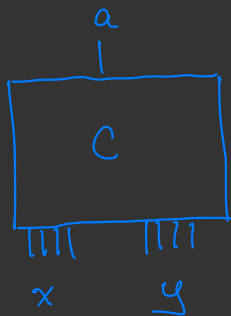
- If not, what is the shortest length of a solution to the problem? What about complexity notions?

A motivation for complexity of sols. vs problems.

Thm (Impagliazzo - Wigderson) unless  $NEXP \subseteq \Sigma_2 \in PH_1$ ,

Succinct-3-coloring (NEXP-complete) does not have succinct solutions!

Succinct-3-coloring:



Input:  $\langle C \rangle \leftarrow$  circuit description

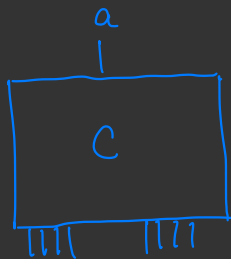
$G =$  graph implicitly defined by  $C$ .

edge  $x \sim y \iff C(x, y) = 1$ .

Goal: Decide if  $G$  is 3-colorable.

A motivation for complexity of sols. vs problems.

Succinct-3-coloring: Input:  $\langle C \rangle \leftarrow$  circuit description



$x, y \in \{0,1\}^n$

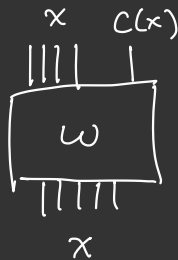
$G =$  graph implicitly defined by  $C$ .

edge  $x \sim y \iff C(x,y) = 1$ .

Goal: Decide if  $G$  is 3-colorable.


Say S3COL instance  $\langle C \rangle$  has a succinct sol. if  $\exists$

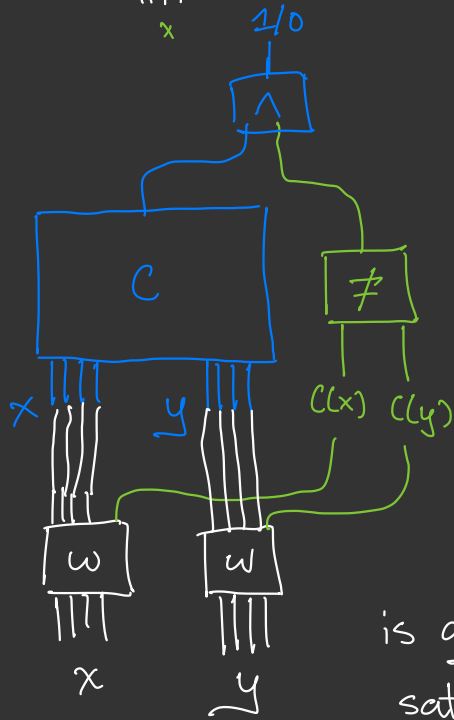
poly sized ckt



outputting coloring  $C(x)$  from  
optimal coloring.

Thm 3BCOL doesn't have succinct sols.


PF If  $\exists$   then

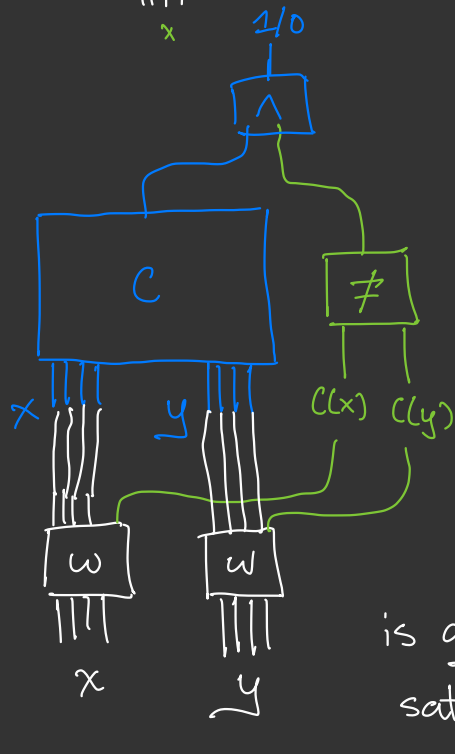


output 1 if

$x \sim y$  AND  $c(x) \neq c(y)$ .

Thm S3COL doesn't have succinct sols.

PF If  $\exists$   then



is always satisfiable.

Then  $\langle C \rangle \in \text{S3COL}$  iff

$\exists \langle w \rangle$  s.t. BIG-CKT is always satisfiable.

i.e.  $\exists w$  s.t.  $\forall x, y$

$$B(x, y, w) = 1$$

$$\Rightarrow \text{NEXP} \subseteq \Sigma_2 \subseteq \text{PH.}$$

Why is this classical CS textbook pt important?

It provides a clear separation between the description complexity of sols. and questions.

Notice, that even with a succinct description of S3COL we would not expect to check the problem in sub-exponential time.

$$P \subseteq NP \subseteq \Sigma_2 \subseteq PH \subseteq \dots \subseteq NEXP$$

Exponential time classes

Instead, description complexity yields a speedup among these large complexity classes that all take exponential time.

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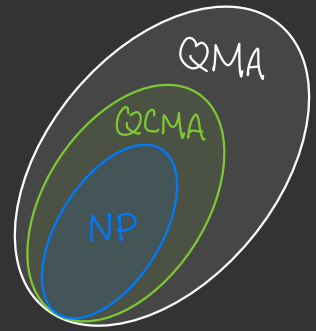
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Today's talk: How should we define description complexity for quantum problems and what is known?



# Non-deterministic quantum computation

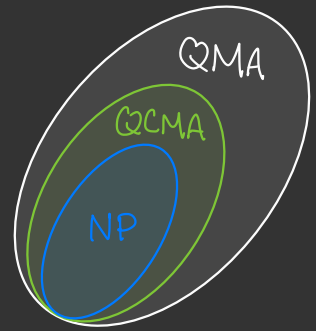
$QMA \stackrel{?}{=} QCMA$ : Do all "classically describable" quantum questions have "classically describable" solutions?



Note: both cases still speculate the problem is exp-hard for BPP (or BQP). It's a matter of description.

# Non-deterministic quantum computation

$QMA \stackrel{?}{=} QCMA$ : Do all "classically describable" quantum questions have "classically describable" solutions?



Note: both cases still speculate the problem is exp-hard for BPP (or BQP). It's a matter of description.

If  $QCMA \neq QMA$ , what complexity class captures the classical complexity of solutions to QMA problems?

## Search-to-decisions:

How much harder is finding a solution than deciding if one exists?

For the class NP, it's equally hard...

$$\exists x_2 \dots x_n, \psi(0, x_2, \dots, x_n) = 1$$

yes  
set  $y_1 = 0$  ↙

no  
set  $y_1 = 1$  ↘

$$\exists x_3 \dots x_n, \psi(y_1, 0, x_3, \dots, x_n) = 1$$

yes  
set  $y_2 = 0$  ↙

no  
set  $y_2 = 1$  ↘

⋮

End of process,  
 $y_1 \dots y_n$  forms a sol.  
to  $\psi$ .

# What about quantum search-to-decision?

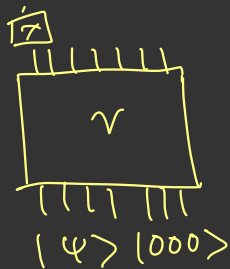
First What does search-to-decision mean in this context?

Issues: 1. QMA is a promise class.

2. The solution might depend on the verifier.

## Search QMA def:

Given a canonical QMA problem described as a verifier  $V$

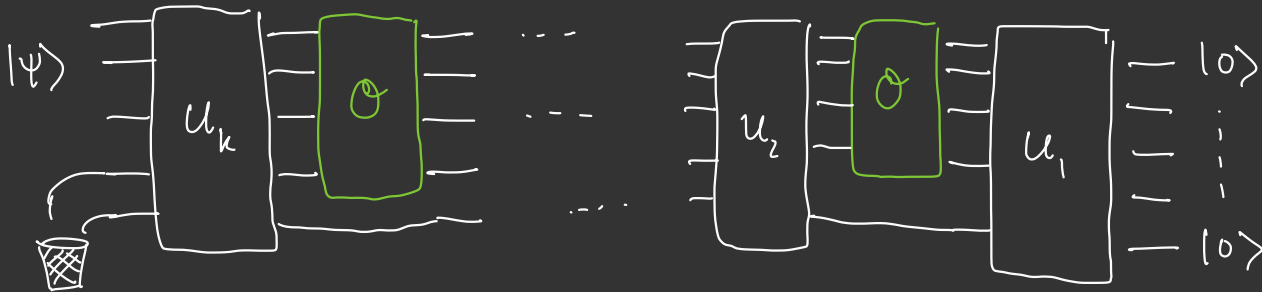


output a state  $|\psi\rangle$  which that verifier will

accept with prob.  $\frac{2}{3}$ .

# QMA search-to-decision reductions

Input: Verifier  $\mathcal{V}$ .



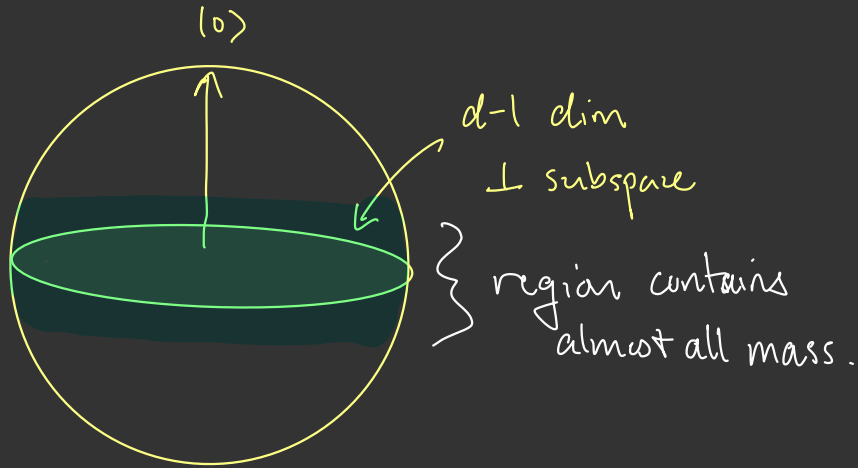
Circuit with oracle gates  $\mathcal{O}$  accessed in superposition

$$\mathcal{O}(x) = \begin{cases} 1 & \text{if } x \text{ encodes a YES QMA question} \\ 0 & \text{if } x \text{ encodes a NO QMA question} \\ \text{either} & \text{if } x \text{ encodes an invalid QMA question} \end{cases}$$

Goal: Output  $|\Psi\rangle$  accepted w pr  $\frac{2}{3}$  by  $\mathcal{V}$ .

## Difficulties to overcome

There is no good way to binary search over the Hilbert space.



Trying to find  $|\psi\rangle$  by a seq. of projectors is a no-go path.

"entanglement destroying"

(also why ground-space dim counting seems hard).

## Thm (Aaronson/Folklore)

$\exists$  a  $2n+1$  query algorithm for generating any state  $|\psi\rangle$  up to  $\exp(-n)$  accuracy.

(When applied to QMA sols., oracle complexity = PP.)

## Our theorem Thm ( $1 \leq n \leq 2^1$ )

$\exists$  a 1-query PP algorithm for generating the sol. to QMA problems.

(We also have extensions to general states).

# Crucial intuitions

① Building all states is unnecessarily powerful.

By counting, there are only  $2^{\text{poly}(n)}$  QMA problems  $\ll \exp(-n)$  net over  $\mathcal{H}$ .  
real vs. imaginarity

② Since QMA states are verifiable, ~~signs~~ of amplitudes don't matter.  $|\psi\rangle = \sum_x \alpha_x |x\rangle$ , then  $\exists |\phi\rangle = \sum_x \beta_x |x\rangle$   $\beta_x \in \mathbb{R}$  s.t.  $|\langle \phi | \psi \rangle| \geq \text{constant}$ .

③ If  $|\psi\rangle$  is Haar-random, then the amplitudes concentrate around  $\frac{1}{\sqrt{2^n}}$ .  $\mathbb{E} |\langle x | \psi \rangle|^2 = \frac{1}{2^n}$ ,  $\mathbb{E} |\langle x | \psi \rangle|^4 = \frac{2}{2^n(2^n + 1)}$ .  
 $|\psi\rangle \sim \text{Haar}$ ,  $|\psi\rangle \sim \text{Haar}$



Interlude: Phase states.  $f: \{0,1\}^n \rightarrow \{0,1\}$

$$|\Psi_f\rangle = \frac{1}{\sqrt{2^n}} \sum_x (-1)^{f(x)} |x\rangle = \text{circuit diagram}$$

The diagram shows a box labeled 'f' with multiple input lines on the left and three output lines on the right. Each output line passes through a box labeled 'F' (representing a Hadamard gate) before reaching a state label: the top line is labeled  $|0\rangle$ , the middle line is labeled  $|0\rangle$ , and the bottom line is labeled  $|0\rangle$ .

For any vector  $|v\rangle \in \mathbb{R}^{2^n}$ , best phase state approx  $|v\rangle$  is with  $f(x) = \text{sgn}(\langle x|v\rangle)$ .

$$\Rightarrow \langle \Psi_f | v \rangle = \frac{\| |v\rangle \|_1}{\sqrt{2^n}}$$

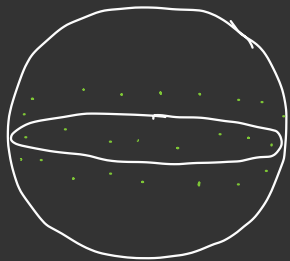
Lem  $\Pr_{|v\rangle \sim \text{Haor}} \left[ \frac{\| |v\rangle \|_1}{\sqrt{2^n}} < \frac{\sqrt{\alpha}}{2} \right] < \alpha.$

## Phase states (cont.)

Lem  $\Pr_{|\nu\rangle \sim \text{Haor}} \left[ \frac{\|\nu\rangle\|_1}{\sqrt{2^n}} < \frac{\sqrt{\alpha}}{2} \right] < \alpha.$

$\Pr_{\text{fig}} \left[ |\langle \Psi_A | \Psi_g \rangle| > \delta \right] \leq 2 \exp\left(-\frac{\delta^2 \cdot 2^n}{3}\right)$  Chernoff bound.

In short, phase states form an effective net for the Hilbert space under the Haar measure.



goal: show PP fn  $f$  s.t.

$|\Psi_A\rangle$  approximates QMA sol.

## Small issues to handle

① Sol.  $|\tau\rangle$  may not be approximable by phase states.

But for Clifford  $C$ ,  $C^\dagger H C$  will be whp.

Then can rotate phase states by  $C^\dagger$  to recover.

② To define fn  $f(x) = \text{sgn}(\text{Tr}(\langle x | \tau \rangle))$  we need

$|\tau\rangle$ . But,

$$|\tau\rangle \propto (\mathbb{1} - H)^{\text{poly}(n)} \underbrace{D|0^n\rangle}_{\text{random clifford state}}$$

$$f(x) = \text{sgn}(\text{Tr}(\langle x | C^\dagger (\mathbb{1} - H)^P D |0^n\rangle)).$$

Thm 1 query PP alg which outputs a state  $|\psi\rangle$

s.t.  $|\langle\psi|\tau\rangle|^2 \geq 2^{-10}$  whp.

- Can add phase estimation to either output  $|\tau\rangle \pm \frac{1}{\text{poly}(n)}$

w pr  $2^{-10}$ .

- Algorithm is parallelizable with still one query

to boost success prob. to  $1 - \frac{1}{\text{poly}(n)}$ .

Is this the best we can do?

Oracle no-go result for QMA-search to QMA-decision  
reduction.

Thm ( $\text{INN}^{21}$ )

QMA<sup>0</sup> search problem with no QMA<sup>0</sup> decision oracle alg.

$O = \mathbb{1} - 2|\psi_f\rangle\langle\psi_f|$  where  $|\psi_f\rangle$  is a phase state

OR  $O = \mathbb{1}$ . Problem: Decide which scenario.

Idea All sols. accepted w pr  $\geq \frac{2}{3}$ , have large support

on  $|\psi_f\rangle$ .

## Oracle no-go (cont.)

PF sketch: (1) Assume  $\exists$  alg  $A^{\text{QMA}^{\Theta}, \Theta}$  that produces  $|\Psi_f\rangle$ .

(2) Show that when run on  $\Theta' = \mathbb{1} - |\Psi_g\rangle\langle\Psi_g|$ ,  
alg's step-by-step behaviour is similar (hybrid alg).

(3) Argue whp should output nearly  $\perp$  states  
and yet cannot by hybrid alg.

# Consequences

QMA sols. can be described by phase states corresponding to PP fns ( $2^{\text{poly}(n)}$  PP fns and  $2^{\text{poly}(n)}$  QMA problems) vs.  $2^{2^n}$  phase states in general

∄ any hope for search-to-decision reductions for generic phase states (which are sols. to QMA<sup>0</sup> problems)

Due to similarity of oracle separating QCMA/QMA, we suspect some oracles show S-to-D No-go's.

## Final thoughts before I finish

- ① Devote more research to understanding descriptions of  $q$ . states. Not the same as decision problems!
- ② Simpler descriptions lead to decision problem speedups.
- ③ Big open questions are
  - ③a  $QCMA \stackrel{?}{=} QMA$
  - ③b Is description complexity robust to small perturbations? i.e. extensions of NLTS theorem.